

No. 9: Application of the Euler-Maclaurin summation to log-differentials of $M(t) = |\zeta(\frac{1}{2} + it)|$

Towards a Proof of the Riemann Hypothesis

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Abstract

This is our first report on numerical evaluation methods for the Riemann zeta function and its related functions. After reviewing the classical Euler-Maclaurin summation, we follow Edwards [2] in the use of this summation formula to evaluate the Riemann zeta function. Then, we extend this method to the evaluation of log-differentials of the modulus $M(t) = |\zeta(\frac{1}{2} + it)|$. We are interested in the function $M(t)$, since it is closely related to the function $\Xi(t) = \zeta(\frac{1}{2} + it)$ and Hardy's Z-function, which we have investigated in our previous reports [5, 6, 7, 8].

Key words: Euler-Maclaurin summation, Bernoulli polynomials, Bernoulli numbers, Riemann zeta function $\zeta(s)$, Modulus function $M(t) = |\zeta(\frac{1}{2} + it)|$, Function $\Xi(t) = \xi(\frac{1}{2} + it)$, Z-function $Z(t)$, Riemann-Siegel theta function $\vartheta(t)$, log-differentials.

1 Introduction: Euler-Maclaurin summation, Bernoulli polynomials, and Bernoulli numbers

The material that we present in this section draws heavily upon Edwards [2], Chapter 6, who gives a comprehensive account of this topic. For a continuously differentiable function $f(x)$ defined for $x \in [M, N]$, we can show that the following formula holds ([2] p. 99, (3)):

$$\sum_{n=M}^N f(n) = \int_M^N f(x) dx + \frac{1}{2}[f(M) + f(N)] + \int_M^N (x - [x] - \frac{1}{2})f'(x) dx, \quad (1)$$

where $[x]$ is the floor function, i.e., $[x]$ is the largest integer that does not exceed x . The third term can be manipulated by using the periodic property of $x - [x] - \frac{1}{2}$, and by using integration by parts:

$$\begin{aligned} \int_M^N (x - [x] - \frac{1}{2})f'(x) dx &= \sum_{n=M}^{N-1} \int_n^{n+1} (x - [x] - \frac{1}{2})f'(x) dx = \sum_{n=M}^{N-1} \int_0^1 (t - \frac{1}{2})f'(n+t) dt \\ &= \sum_{n=M}^{N-1} \left[(t - \frac{1}{2})f(n+t) \Big|_0^1 - \int_0^1 f(n+t) dt \right] \\ &= \sum_{n=M}^{N-1} \left[\frac{1}{2}f(n+1) + \frac{1}{2}f(n) \right] - \sum_{n=M}^{N-1} \int_0^1 f(n+t) dt \\ &= \frac{1}{2}f(M) + f(M+1) + \cdots + f(N-1) + \frac{1}{2}f(N) - \int_M^N f(x) dx, \end{aligned} \quad (2)$$

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which indeed proves the formula (1). This formula can be transformed into a useful form by further use of integration by parts, and by adopting the Bernoulli polynomials, as will be shown below.

The n th Bernoulli polynomials $B_n(x)$ is the unique polynomial of degree n with the property

$$\int_x^{x+1} B_n(t) dt = x^n. \quad (3)$$

By differentiating the above, we obtain

$$B_n(x+1) - B_n(x) = \int_x^{x+1} B'_n(t) dt = nx^{n-1}. \quad (4)$$

The last equality implies that $\frac{B'_n(x)}{n}$ satisfies the definition of $B_{n-1}(x)$, which is the integrand in (3), with n replaced by $n-1$. Hence,

$$B'_n(x) = nB_{n-1}(x). \quad (5)$$

The first four Bernoulli polynomials are given as

$$B_0(x) = 1, \quad B_1(x) = x - \frac{1}{2}, \quad B_2(x) = x^2 - x + \frac{1}{6}, \quad B_3(x) = x^3 - \frac{3}{2}x^2 + \frac{1}{2}x. \quad (6)$$

Using these polynomials, we have from the last expression of the first line of (2):

$$\int_M^N (x - [x] - \frac{1}{2})f'(x) dx = \sum_{n=M}^{N-1} \int_0^1 B_1(t)f'(n+t) dt = \sum_{n=M}^{N-1} \left[\frac{B_2(t)}{2}f'(n+t) \Big|_0^1 - \frac{1}{2} \int_0^1 B_2(t)f''(n+t) dt \right], \quad (7)$$

where the relation (5) with $n=2$ is used. By noting $B_2(0) = B_2(1) (= \frac{1}{6})$, the first term of RHS in the above can be written as

$$\sum_{n=M}^{N-1} \frac{B_2(t)}{2}f'(n+t) \Big|_0^1 = \frac{B_2(0)}{2} \sum_{n=M}^{N-1} [f'(n+1) - f'(n)] = \frac{B_2(0)}{2} f'(x) \Big|_M^N. \quad (8)$$

The last term in (7) can be written, by defining the periodic function $B_2(x - [x]) = \bar{B}_2(x)$, as

$$\begin{aligned} \sum_{n=M}^{N-1} \frac{1}{2} \int_0^1 B_2(t)f''(n+t) dt &= \frac{1}{2} \int_M^N B_2(x - [x])f''(x) dx = \frac{1}{2} \int_M^N \bar{B}_2(x)f''(x) dx = \frac{1}{2 \cdot 3} \int_M^N \bar{B}'_3(x)f''(x) dx \\ &= \frac{1}{2 \cdot 3} \bar{B}_3(x)f''(x) \Big|_M^N - \frac{1}{2 \cdot 3} \int_M^N \bar{B}_3(x)f'''(x) dx. \end{aligned} \quad (9)$$

The identity $B_3(0) = B_3(1) = 0$ implies $\bar{B}_3(n) = 0$ for any integer n , making the first term in RHS of (9) disappear. Thus, (7) can be simplified as

$$\int_M^N (x - [x] - \frac{1}{2})f'(x) dx = \frac{B_2(0)}{2} f'(x) \Big|_M^N + \frac{1}{2 \cdot 3} \int_M^N \bar{B}_3(x)f'''(x) dx. \quad (10)$$

By carrying integration by parts further, we obtain

$$\begin{aligned} \sum_{n=M}^N f(n) &= \int_M^N f(x) dx + \frac{1}{2}[f(M) + f(N)] + \frac{B_2(0)}{2} f'(x) \Big|_M^N - \frac{B_3(0)}{2 \cdot 3} f''(x) \Big|_M^N + \cdots + (-1)^k \frac{B_k(0)}{k!} f^{(k-1)}(x) \Big|_M^N \\ &\quad + (-1)^{k+1} \frac{1}{k!} \int_M^N \bar{B}_k(x)f^{(k)}(x) dx. \end{aligned} \quad (11)$$

The Bernoulli polynomial $B_n(x)$ reduces to the Bernoulli number B_n at $x = 0$. The Bernoulli numbers B_n are defined as the coefficient of the Taylor series expansion of the generating function $\frac{te^t}{e^t-1}$ (see e.g. [1] Chapter 1,[2], p. 11, [4], p.7)

$$\frac{te^t}{e^t-1} = \sum_{n=0}^{\infty} B_n \frac{x^n}{n!}. \quad (12)$$

Another definition of the Bernoulli numbers is by the following recursion:

$$\sum_{k=0}^{n-1} \binom{n}{k} B_k = n, \quad n = 1, 2, 3, \dots \quad (13)$$

The first several B_n are

$$B_0 = 1, \quad B_1 = -\frac{1}{2}, \quad B_2 = \frac{1}{6}, \quad B_3 = 0, \quad B_4 = -\frac{1}{30}. \quad (14)$$

Furthermore, the Bernoulli numbers B_n with odd n are all zero, except for B_1 , i.e.,

$$B_{2k+1} = 0 \quad \text{for all } k \geq 1. \quad (15)$$

Substituting these properties, we finally obtain what is known as the *Euler-Maclaurin summation formula*:

$$\sum_{n=M}^N f(n) = \int_M^N f(x) dx + \frac{1}{2}[f(M) + f(N)] + \frac{B_2}{2} f'(x)|_M^N + \frac{B_4}{4!} f'''(x)|_M^N + \dots + \frac{B_{2\nu}}{(2\nu)!} f^{(2\nu-1)}(x)|_M^N + R_{2\nu}, \quad (16)$$

where the error term can be expressed by setting $k = 2\nu$ in (11)

$$R_{2\nu} = -\frac{1}{(2\nu)!} \int_M^N \bar{B}_{2\nu}(x) f^{(2\nu)}(x) dx. \quad (17)$$

Since the term containing $B_{2\nu+1}$ can be ignored, we may alternatively set $k = 2\nu + 1$, yielding

$$R_{2\nu} = R_{2\nu+1} = \frac{1}{(2\nu+1)!} \int_M^N \bar{B}_{2\nu+1}(x) f^{(2\nu+1)}(x) dx. \quad (18)$$

A simple estimate of the error term $|R_{2\nu}|$ is obtained if $f^{(2\nu+1)}(x)$ is monotone decreasing, by using the fact that $\bar{B}_{2\nu+1}(x)$ alternates in sign as $\bar{B}_3(x)$ does. Edwards (cf. [2], p. 105) states “It is a general rule of thumb in applying the Euler-Maclaurin summation formula that as long as the terms are decreasing rapidly in size, the bulk of the errors is in the first term omitted.”

2 Evaluation of $\zeta(s)$ by Euler-Maclaurin summation

We continue to follow the exposition by Edwards [2], Section 6.4. Instead of directly applying the Euler-Maclaurin summation method to $\zeta(s) = \sum_{n=1}^{\infty} n^{-s}$, we consider summing $f(x) = x^{-s}$ from $x = N$ to infinity:

$$\begin{aligned} \sum_{n=N}^{\infty} n^{-s} &= \zeta(s) - \sum_{n=1}^{N-1} n^{-s} = \int_N^{\infty} f(x) dx + \frac{1}{2}[f(N) + f(\infty)] + \int_N^{\infty} \bar{B}_1(x) f'(x) dx \\ &= \int_N^{\infty} x^{-s} dx + \frac{1}{2}N^{-s} - s \int_N^{\infty} \bar{B}_1(x) x^{-s-1} dx \\ &= \frac{N^{1-s}}{s-1} + \frac{1}{2}N^{-s} + \frac{B_2}{2} s N^{-s-1} + \dots + \frac{B_{2\nu}}{(2\nu)!} s(s+1) \dots (s+2\nu-2) N^{-s-2\nu+1} + R_{2\nu}, \end{aligned} \quad (19)$$

where

$$R_{2\nu} = \frac{1}{(2\nu+1)} \int_N^\infty \bar{B}_{2\nu+1}(x) f^{(2\nu+1)}(x) dx = -\frac{s(s+1)\cdots(s+2\nu)}{(2\nu+1)!} \int_N^\infty \bar{B}_{2\nu+1}(x) x^{-s-2\nu-1} dx. \quad (20)$$

By rearranging (19), we have the Euler-Maclaurin sum representation for $\zeta(s)$:

$$\zeta(s) = \sum_{n=1}^{N-1} n^{-s} + \frac{N^{1-s}}{s-1} + \frac{1}{2}N^{-s} + \frac{B_2}{2}sN^{-s-1} + \cdots + \frac{B_{2\nu}}{(2\nu)!}s(s+1)\cdots(s+2\nu-2)N^{-s-2\nu+1} + R_{2\nu}. \quad (21)$$

If N is about the same size as $|s|$, the terms of the series (21) decreases quite rapidly, so we expect the reminder $R_{2\nu}$ to be quite small. We wish to obtain an upper bound of its magnitude.

$$\begin{aligned} |R_{2\nu}| &= \left| \frac{s(s+1)\cdots(s+2\nu)}{(2\nu+1)!} \int_N^\infty \bar{B}_{2\nu-1}(x) x^{-s-2\nu+1} dx \right| \\ &= \left| \frac{s(s+1)\cdots(s+2\nu-1)}{(2\nu)!} \int_N^\infty [B_{2\nu+2} - \bar{B}_{2\nu+2}(x)] x^{-s-2\nu-2} dx \right| \\ &\leq \frac{|s(s+1)\cdots(s+2\nu-1)|}{(2\nu)!} \int_N^\infty |B_{2\nu+2} - \bar{B}_{2\nu+2}(x)| x^{-\sigma-2\nu-2} dx. \end{aligned} \quad (22)$$

where the second line of the above was obtained by noting that $[B_{2\nu+2} - \bar{B}_{2\nu+2}(x)]' = -(2\nu+2)\bar{B}_{2\nu+1}(x)$, and $B_{2\nu+2} - \bar{B}_{2\nu+2}(N) = 0$ after integration by parts, and where $\sigma = \Re(s)$ in the last line.

It can be shown that $B_{2\nu+2} - \bar{B}_{2\nu+2}(x)$ is zero at $x = 0$ and $x = 1$ and has its extremum at $x = \frac{1}{2}$ and this is the only extremum in $0 < x < 1$ (see e.g., [2] p. 112). This implies that $B_{2\nu+2} - \bar{B}_{2\nu+2}(x)$ never changes its sign, and has the same sign as $B_{2\nu+2}$ (because $B_{2\nu+2}(x)$ takes zero at $x = 0$ and $x = 1$). If we recall Euler's formula $\zeta(2\nu) = (-1)^{\nu-1} \frac{(2\pi)^{2\nu}}{2(2\nu)!} B_{2\nu}$ (see e.g. [4], p. 7, Eqn. (42)), we readily find that this sign is $(-1)^\nu$, i.e.,

$$|B_{2\nu+2} - \bar{B}_{2\nu+2}(x)| = (-1)^\nu [B_{2\nu+2} - \bar{B}_{2\nu+2}(x)]. \quad (23)$$

Thus, the integration in the last expression of (22) becomes

$$\begin{aligned} \int_N^\infty |B_{2\nu+2} - \bar{B}_{2\nu+2}(x)| x^{-\sigma-2\nu-2} dx &= \int_N^\infty (-1)^\nu [B_{2\nu+2} - \bar{B}_{2\nu+2}(x)] x^{-\sigma-2\nu-2} dx \\ &= |B_{2\nu+2}| \int_N^\infty x^{-\sigma-2\nu-2} dx - \int_N^\infty (-1)^\nu \bar{B}_{2\nu+2}(x) x^{-\sigma-2\nu-2} dx \\ &= \frac{|B_{2\nu+2}|}{\sigma+2\nu+1} N^{-\sigma-2\nu-1} - \frac{|\sigma+2\nu+2|}{2\nu+3} \int_N^\infty (-1)^\nu \bar{B}_{2\nu+3}(x) x^{-\sigma-2\nu-3} dx \\ &\leq \frac{|B_{2\nu+2}|}{|\sigma+2\nu+1|} N^{-\sigma-2\nu-1}. \end{aligned} \quad (24)$$

By combining (22) and (24), we finally obtain

$$\begin{aligned} |R_{2\nu}| &\leq \frac{|s(s+1)\cdots(s+2\nu+1)B_{2\nu+2}N^{-\sigma-2\nu-1}|}{(2\nu+2)!|\sigma+2\nu+1|} \\ &= \frac{|s+2\nu+1|}{|\sigma+2\nu+1|} \cdot \left| \frac{B_{2\nu+2}}{(2\nu+2)!} s(s+1)\cdots(s+2\nu)N^{-s-2\nu-1} \right|. \end{aligned} \quad (25)$$

Note that the $B_{2\nu+2}$ term is the first term omitted in the series representation (21). In other words, the error term $R_{2\nu} = R_{2\nu+1}$ is at most $\frac{|s+2\nu+1|}{|\sigma+2\nu+1|}$ times the magnitude of the first term omitted. Edwards ascribes this result to Backlund's dissertation in 1916 (see [2], p. 115).

Although the Euler-Maclaurin summation formula (21) was derived from the definition $\zeta(s) = \sum_{n=1}^\infty n^{-s}$ which is defined only for $\Re(s) > 1$, the formula is valid throughout the half-plane $\Re(s+2\nu+1) > 1$, in which the integral for $R_{2\nu}$ converges. Since ν is arbitrary, this provides an alternative proof for analytical continuation of $\zeta(s)$ throughout the entire s -plane, except at $s = 0$, which is a simple pole.

3 Evaluation of log differentials of $M(t) = |\zeta(\frac{1}{2} + it)|$

The materials we present in this section is new to the best of this author's knowledge. If we write

$$\zeta(\frac{1}{2} + it) = \sum_{n=1}^{\infty} n^{-\frac{1}{2}-it} = \sum_{n=1}^{\infty} \frac{\cos(t \log n)}{\sqrt{n}} - i \sum_{n=1}^{\infty} \frac{\sin(t \log n)}{\sqrt{n}} = A(t) - iB(t), \quad (26)$$

its modulus $M(t)$ is given by

$$M(t) = |\zeta(\frac{1}{2} + it)| = \sqrt{A^2(t) + B^2(t)}, \quad (27)$$

and its first and second differentials are

$$M'(t) = \frac{A(t)A'(t) + B(t)B'(t)}{M(t)}, \quad \text{and} \quad M''(t) = \frac{A(t)A''(t) + A'^2(t) + B(t)B''(t) + B'^2(t) - M'^2(t)}{M(t)}, \quad (28)$$

respectively. Thus, the log-differential of the modulus is

$$\frac{d \log M(t)}{dt} = \frac{M'(t)}{M(t)} = \frac{A(t)A'(t) + B(t)B'(t)}{A^2(t) + B^2(t)}. \quad (29)$$

and the second log-differential is

$$\frac{d^2 \log M(t)}{dt^2} = \frac{A(t)A''(t) + A'^2(t) + B(t)B''(t) + B'^2(t)}{A^2(t) + B^2(t)} - 2 \left(\frac{M'(t)}{M(t)} \right)^2. \quad (30)$$

The last expression can be alternatively written as

$$\frac{d^2 \log M(t)}{dt^2} = \frac{M''(t)}{M(t)} - \left(\frac{M'(t)}{M(t)} \right)^2. \quad (31)$$

Recall Hardy' Z-function ([3]. See also [2] p. 119, and [7]) defined by

$$Z(t) = e^{i\vartheta(t)} \zeta(\frac{1}{2} + it), \quad (32)$$

where $\vartheta(t)$ is the Riemann-Siegel theta function:

$$\vartheta(t) = \Im\{\log \Gamma(\frac{1}{4} + \frac{it}{2})\} - \frac{t \log \pi}{2} = \arg\{\zeta(\frac{1}{2} + it)\}. \quad (33)$$

Then, the following identity readily follows from the definition (32):

$$|Z(t)| = M(t). \quad (34)$$

Thus, we have

$$\frac{M'(t)}{M(t)} = \frac{|Z(t)|'}{|Z(t)|} = \frac{Z'(t)}{Z(t)}, \quad (35)$$

where the last equality holds because $Z(t)$ is a real function.

By expanding the right-hand side (RHS) of (32) and setting it equal to LHS, we find

$$Z(t) = A(t) \cos \vartheta(t) + B(t) \sin \vartheta(t), \quad (36)$$

and

$$A(t) \sin \vartheta(t) - B(t) \cos \vartheta(t) = 0, \quad (37)$$

from which we have

$$B(t) = A(t) \tan \vartheta(t). \quad (38)$$

By substituting this into (36), we find

$$Z(t) = \frac{A(t)}{\cos \vartheta(t)} = \frac{B(t)}{\sin \vartheta(t)}. \quad (39)$$

Therefore, we have

$$A(t) = Z(t) \cos \vartheta(t), \quad \text{and} \quad B(t) = Z(t) \sin \vartheta(t), \quad (40)$$

which could have been directly derived from (32).

By differentiating (39), we have

$$Z'(t) = \frac{A'(t)}{\cos \vartheta(t)} + \frac{A(t)\vartheta'(t) \sin \vartheta(t)}{\cos^2 \vartheta(t)} \quad (41)$$

$$= \frac{B'(t)}{\sin \vartheta(t)} - \frac{B(t)\vartheta'(t) \cos \vartheta(t)}{\sin^2 \vartheta(t)}. \quad (42)$$

Thus, by taking the ratio of $Z'(t)$ to $Z(t)$ of (39), we have

$$\frac{Z'(t)}{Z(t)} = \frac{A'(t)}{A(t)} + \vartheta'(t) \tan \vartheta(t) \quad (43)$$

$$= \frac{B'(t)}{B(t)} - \vartheta'(t) \cot \vartheta(t). \quad (44)$$

By differentiating the above, we find

$$\begin{aligned} \frac{d}{dt} \left(\frac{Z'(t)}{Z(t)} \right) &= \frac{d}{dt} \left(\frac{A'(t)}{A(t)} \right) + \frac{d[\vartheta'(t) \tan \vartheta(t)]}{dt} \\ &= \frac{A''(t)}{A(t)} - \left(\frac{A'(t)}{A(t)} \right)^2 + \vartheta''(t) \tan \vartheta(t) + [\vartheta'(t) \sec \vartheta(t)]^2. \end{aligned} \quad (45)$$

Alternatively we have

$$\begin{aligned} \frac{d}{dt} \left(\frac{Z'(t)}{Z(t)} \right) &= \frac{d}{dt} \left(\frac{B'(t)}{B(t)} \right) - \frac{d[\vartheta'(t) \cot \vartheta(t)]}{dt} \\ &= \frac{B''(t)}{B(t)} - \left(\frac{B'(t)}{B(t)} \right)^2 - \vartheta''(t) \cot \vartheta(t) + [\vartheta'(t) \csc \vartheta(t)]^2. \end{aligned} \quad (46)$$

The function $\vartheta(t)$ can have the following Stirling series approximation (see Edwards [2] pp. 119-120):

$$\vartheta(t) = \frac{t}{2} \log \frac{t}{2\pi e} - \frac{\pi}{8} + \frac{1}{48t} + \frac{7}{5760t^3} + \dots \quad (47)$$

4 Evaluation of $\zeta'(s)$

The Euler-Maclaurin summation of $\zeta'(s)$ can be obtained in a manner similar to what we discussed in the preceding section. Since the $n = 1$ term disappear with differentiation, we have

$$\zeta'(s) = \sum_{n=2}^{\infty} \frac{d}{ds} n^{-s} = \sum_{n=2}^{\infty} \frac{d}{ds} e^{-s \log n} = - \sum_{n=2}^{\infty} (\log n) n^{-s}. \quad (48)$$

Thus, by defining $f(x)$ as

$$f(x) = -(\log x)x^{-s} = -(\log x)e^{-s \log x} = a(x)b(x), \quad \text{where } a(x) = \log x, \quad \text{and } b(x) = -x^{-s}. \quad (49)$$

its n -th derivative is given by

$$f^{(n)} = \sum_{k=0}^n \binom{n}{k} a^{(n-k)}(x)b^{(k)}(x), \quad (50)$$

where

$$a^{(j)}(x) = (-1)^{j-1}(j-1)!x^{-j}, \quad j \geq 1, \quad (51)$$

and

$$b^{(k)}(x) = s(s+1)(s+2) \cdots (s+k-1)(-1)^{k-1}x^{-s-k}, \quad k \geq 1. \quad (52)$$

Therefore, by substituting (51) and (52) into (50), we find

$$\begin{aligned} f^{(n)}(x) &= a^{(n)}(x)b(x) + \sum_{k=1}^{n-1} \binom{n}{k} a^{(n-k)}(x)b^{(k)}(x) + a(x)b^{(n)}(x) \\ &= (-1)^n(n-1)!x^{-n-s} + \sum_{k=1}^{n-1} (-1)^n \binom{n}{k} (n-k-1)!s(s+1) \cdots (s+k-1)x^{-n-s} \\ &\quad + (-1)^n s(s+1) \cdots (s+n-1) \log x \cdot x^{-s-n} \\ &= (-1)^n(n-1)!x^{-n-s} \left[1 - s(s+1) \cdots (s+n-1) \log x + \sum_{k=1}^{n-1} \frac{n!}{(n-k)k!} s(s+1) \cdots (s+k-1) \right], \quad n \geq 1. \end{aligned} \quad (53)$$

Thus,

$$f^{(n)}(x) \Big|_N^\infty = (-1)^{n-1}(n-1)!N^{-n-s} \left[1 - s(s+1) \cdots (s+n-1) \log N + \sum_{k=1}^{n-1} \frac{n!}{(n-k)k!} s(s+1) \cdots (s+k-1) \right], \quad n \geq 1. \quad (54)$$

By noting

$$\begin{aligned} \int_N^\infty f(x) dx &= - \int_N^\infty (\log x)x^{-s} dx = - \int_N^\infty \frac{(x^{-s+1})'}{s-1} \log x dx \\ &= \left[\frac{x^{-s+1} \log x}{s-1} \right]_N^\infty + \int_N^\infty \frac{x^{-s+1}}{s-1} \frac{1}{x} dx = \frac{N^{-s+1} \log N}{s-1} + \int_N^\infty \frac{x^{-s}}{s-1} dx \\ &= \frac{N^{-s+1} \log N}{s-1} + \frac{N^{-s+1}}{(s-1)^2} = \frac{N^{-s+1}}{s-1} \left[\log N + \frac{1}{s-1} \right], \end{aligned} \quad (55)$$

and

$$\frac{1}{2}[f(N) + f(\infty)] = -\frac{1}{2}N^{-s} \log N, \quad \text{for } \sigma = \Re(s) > 0, \quad (56)$$

we have the following expression for the Euler-Maclaurin sum series for $\zeta'(s)$:

$$\begin{aligned} \zeta'(s) &= - \sum_{n=2}^{N-1} (\log n)n^{-s} + \frac{N^{-s+1}}{s-1} \left[\log N + \frac{1}{s-1} \right] - \frac{1}{2}N^{-s} \log N + \frac{B_2}{2} f'(x) \Big|_N^\infty + \frac{B_4}{4!} f'''(x) \Big|_N^\infty + \cdots \\ &\quad + \frac{B_{2\nu}}{(2\nu)!} f^{(2\nu-1)}(x) \Big|_N^\infty + R_{2\nu}, \end{aligned} \quad (57)$$

where $f^{(n)}(x) \Big|_N^\infty$ ($n \geq 1$) is given by (54), and the error term is

$$R_{2\nu} = -\frac{1}{(2\nu)!} \int_N^\infty \bar{B}_{2\nu}(x) f^{(2\nu)}(x) dx. \quad (58)$$

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